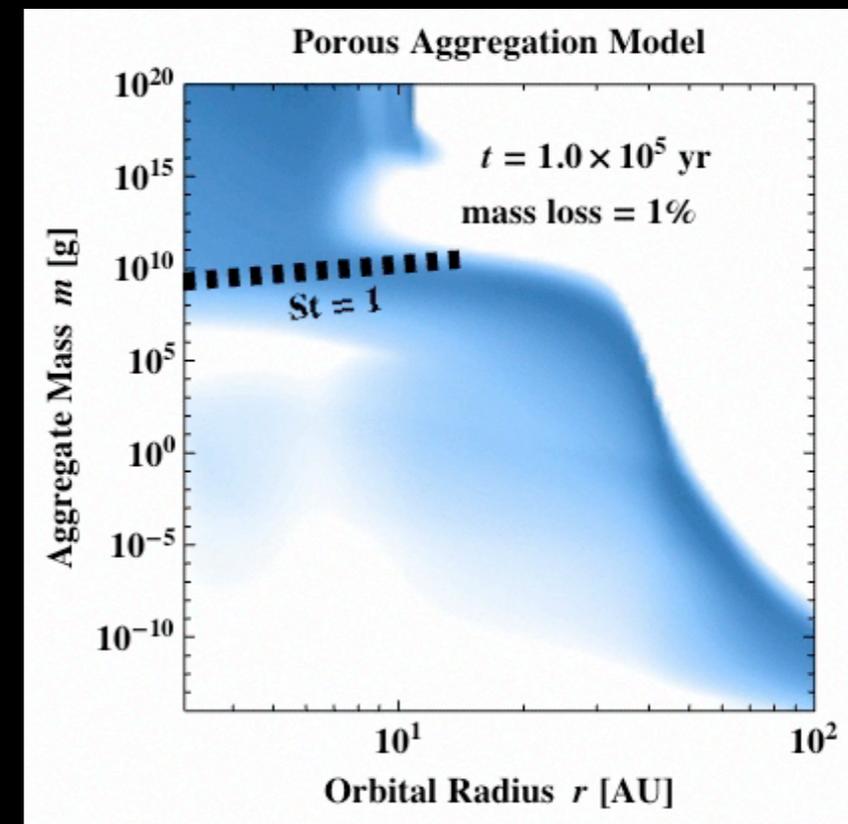


Rapid Coagulation of Porous Dust Aggregates Outside the Snow Line

A Pathway to Successful Icy Planetesimal Formation



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ANIMATION: N-body Collision Experiments by Suyama et al. (2008)

The Radial Drift Barrier against Dust Growth

Headwind \Rightarrow Angular mom. loss
 \Rightarrow Inward drift

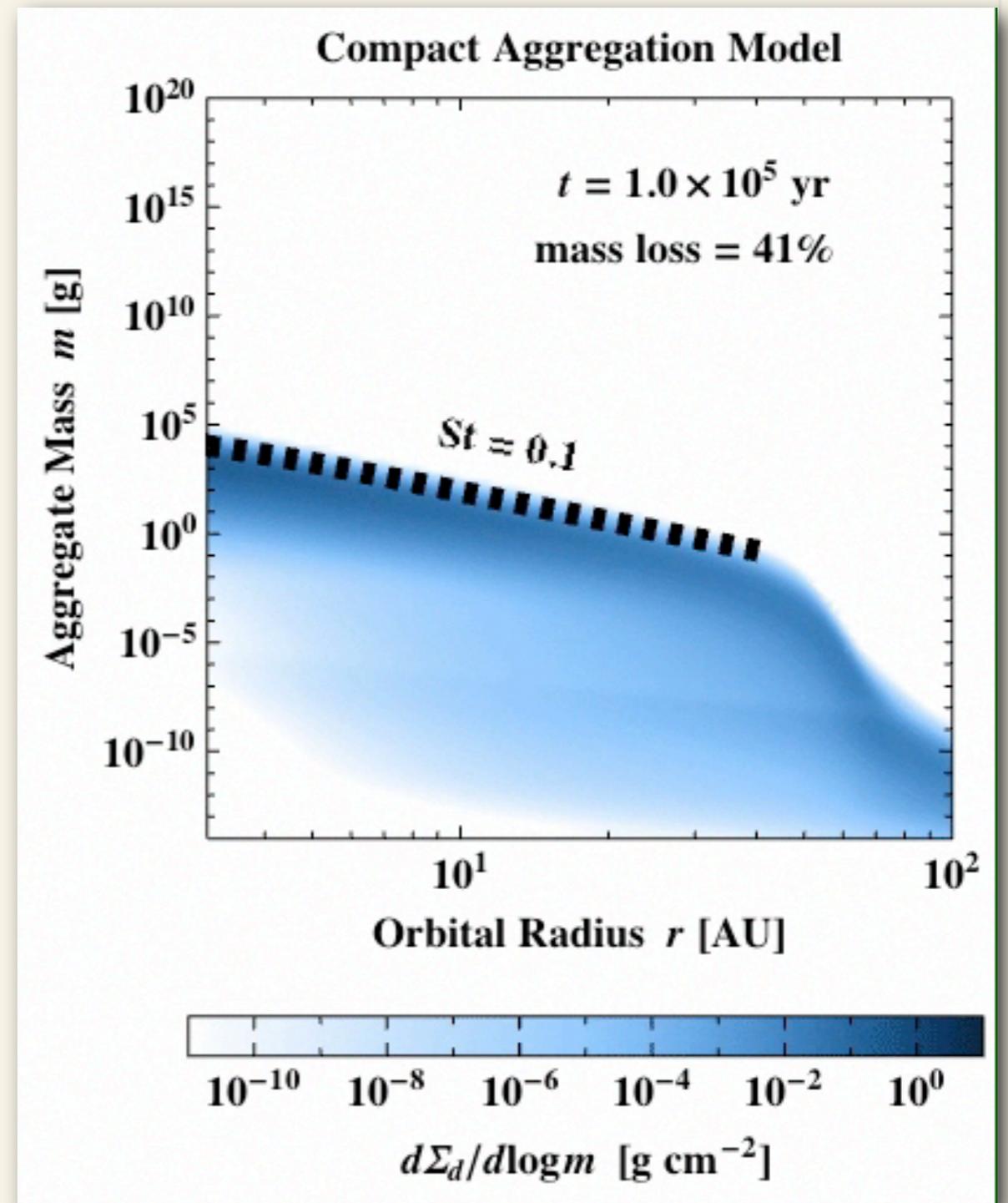
$$v_{\text{drift}} \approx 50 \text{ m s}^{-1} \frac{2\text{St}}{1 + \text{St}^2}$$

$\text{St} = \Omega t_{\text{stop}}$: Stokes number

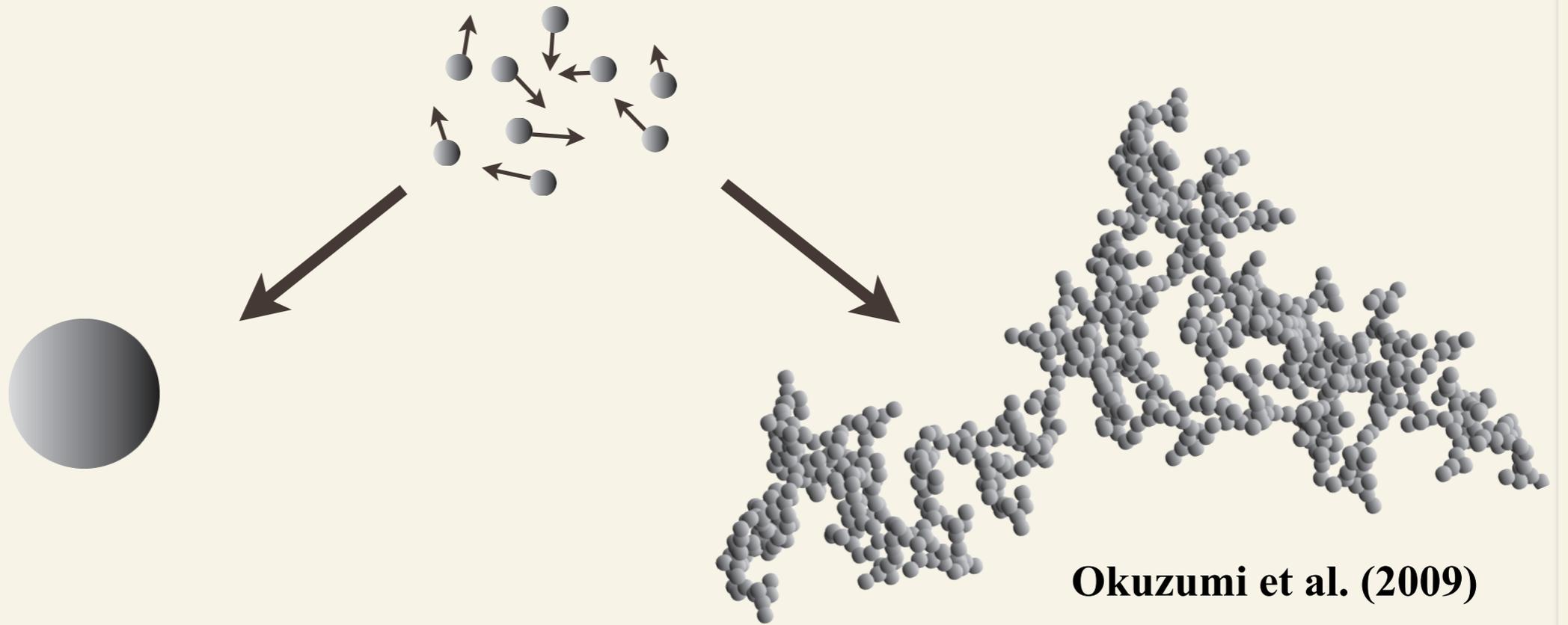
(Adachi et al. 76; Weidenschilling 77)

At $\text{St} \sim 1$, inward drift is faster than collisional growth (Brauer et al. 2008)

\Rightarrow **A barrier against planetesimal formation**



Beyond the Compact Grain Paradigm



Okuzumi et al. (2009)

Compact Sphere

- Filling factor ~ 1
- High St
- The “classic” model

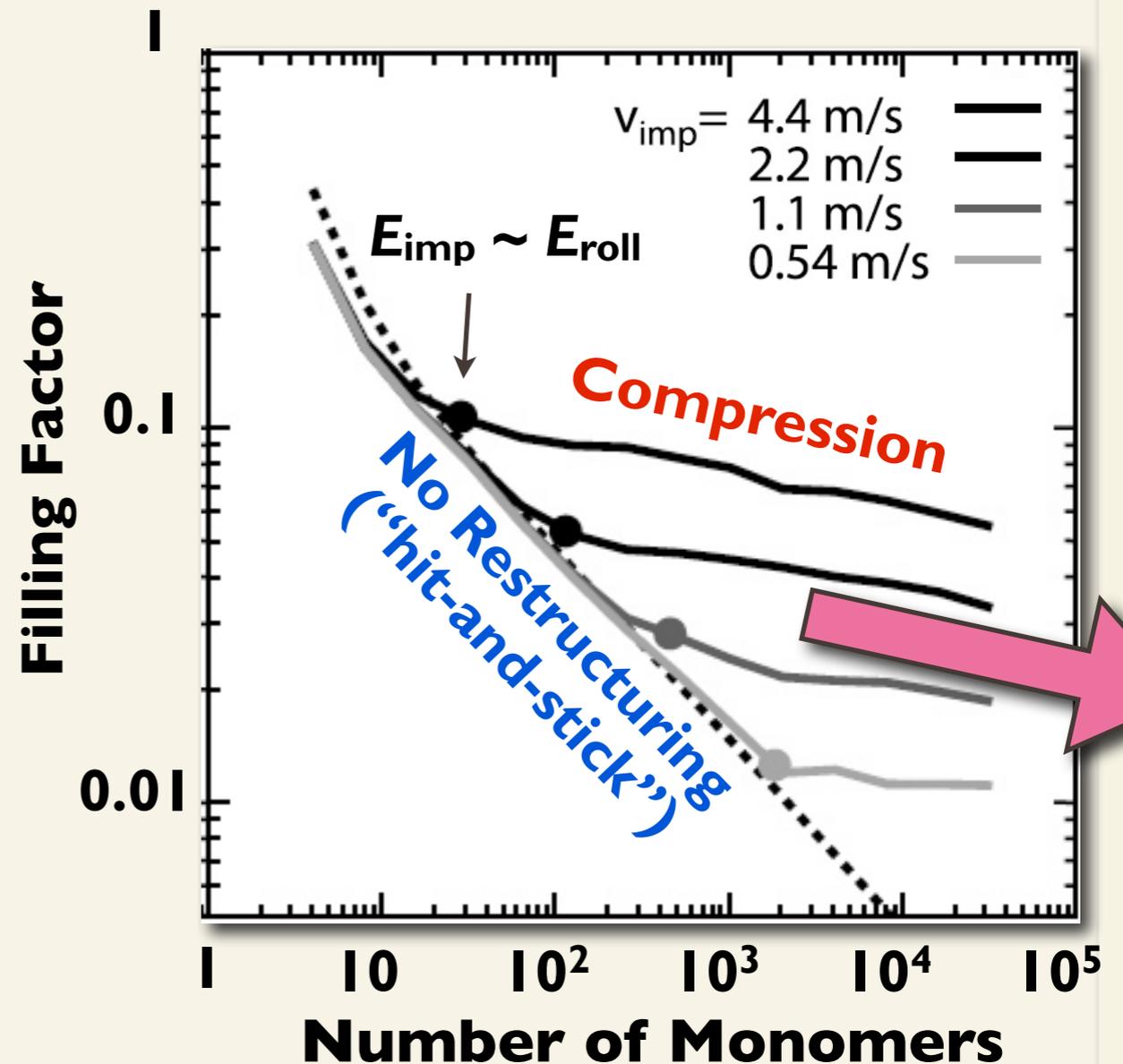
Porous Aggregate

- Filling factor $\ll 1$
- Low St
- Ormel et al. '07; Okuzumi et al. '09; Zsom et al. '10, '11

Collisional Compaction is Inefficient

Suyama, Wada, & Tanaka (2008)

Suyama, Wada, Tanaka, & Okuzumi (2012)



**The filling factor can be $\ll 0.1$
even after the onset of collisional compression!**

Dust Growth Outside the Snow Line

● Coagulation + Radial Drift + Porosity Evolution

- Coagulation + drift: Global bin method (Brauer et al. 2008)
- Porosity: Compression recipe based on N -body experiments (Suyama et al. 2012)

● Outside the Snow Line

- Ice mantle → High sticking efficiency
- Disruption velocity $v_{\text{disr}} \sim 40\text{-}70 \text{ m/s}$ (Wada et al. 2009) \gtrsim Collision velocity

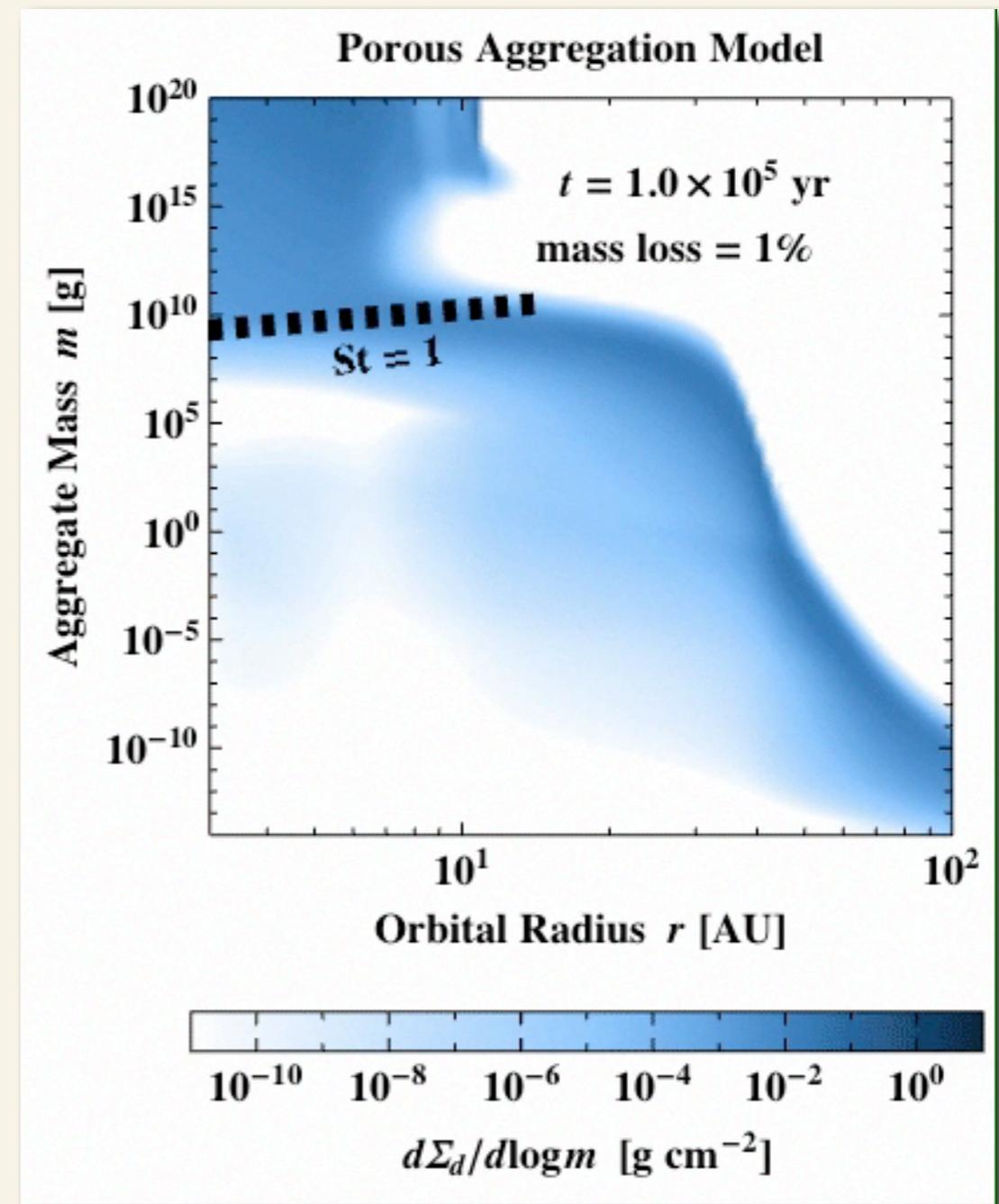
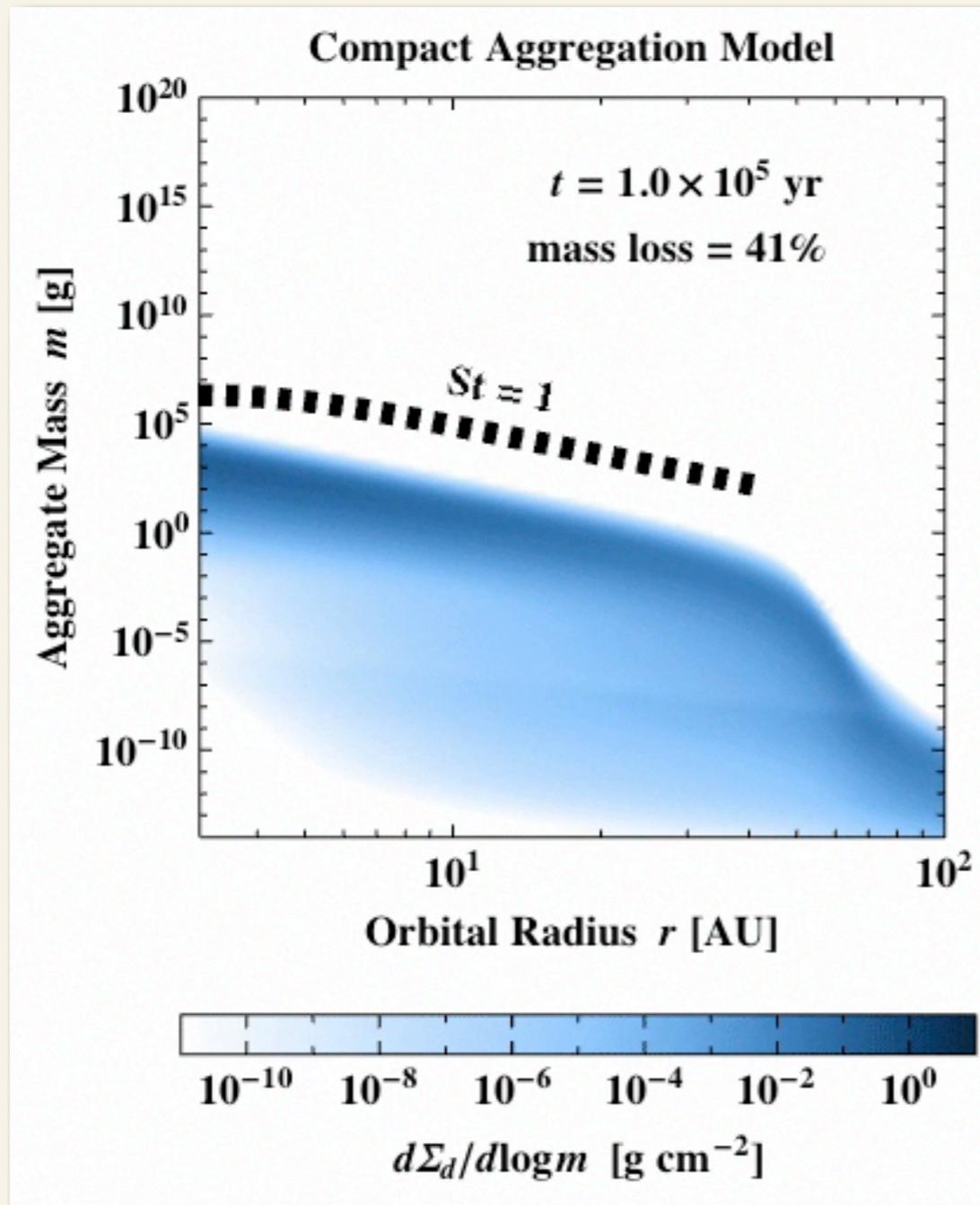
⇒ **We neglect fragmentation**

- Minimum-mass solar nebula (Hayashi et al. 81)
- Turbulence $\alpha = 10^{-3}$
- Dust-to-gas mass ratio = 0.01
- $T = 280(r/\text{AU})^{-1/2} \text{ K} \Rightarrow r_{\text{snow}} \approx 3\text{AU}$

Results: Global Evolution

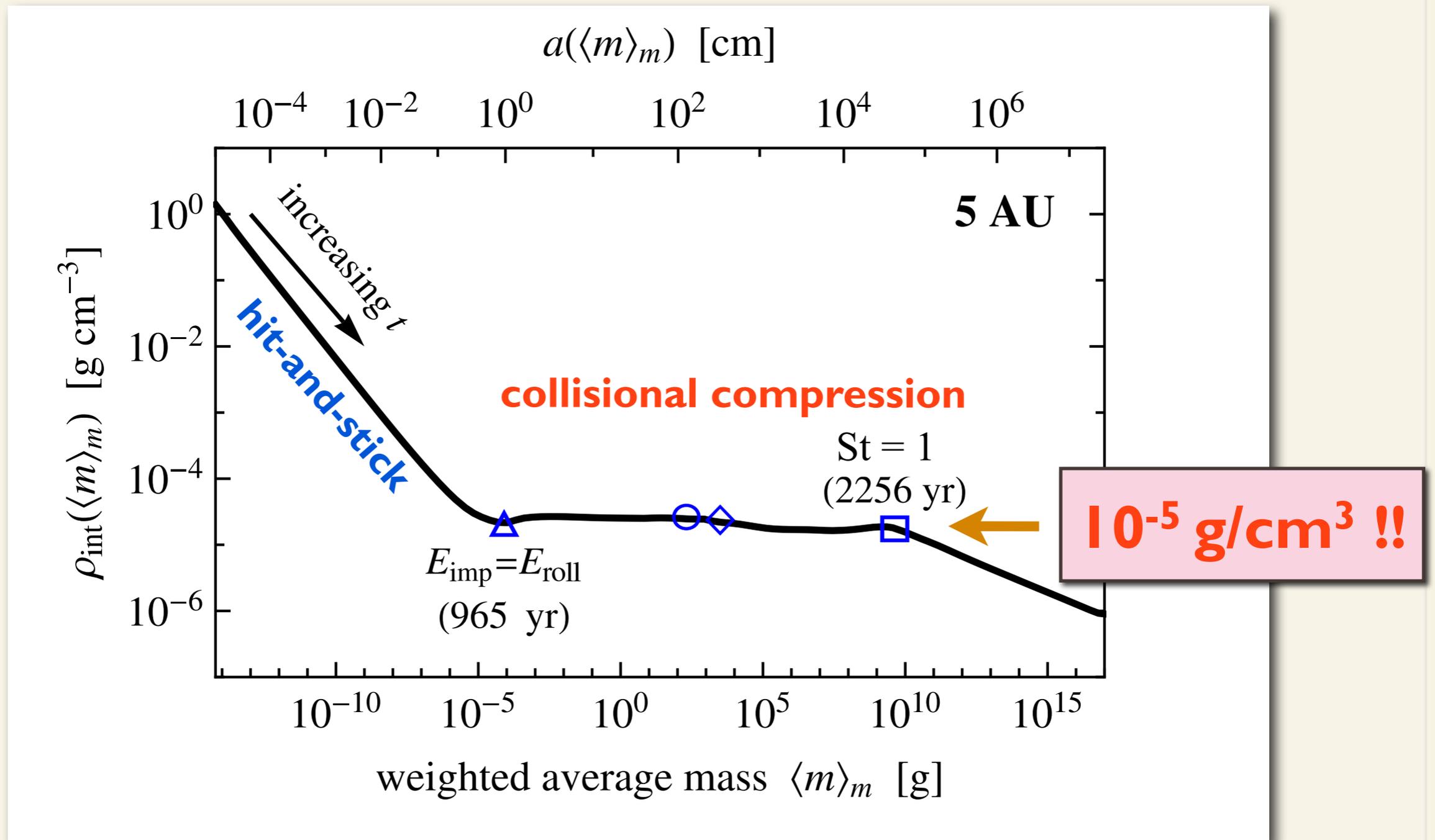
Compact

Porous



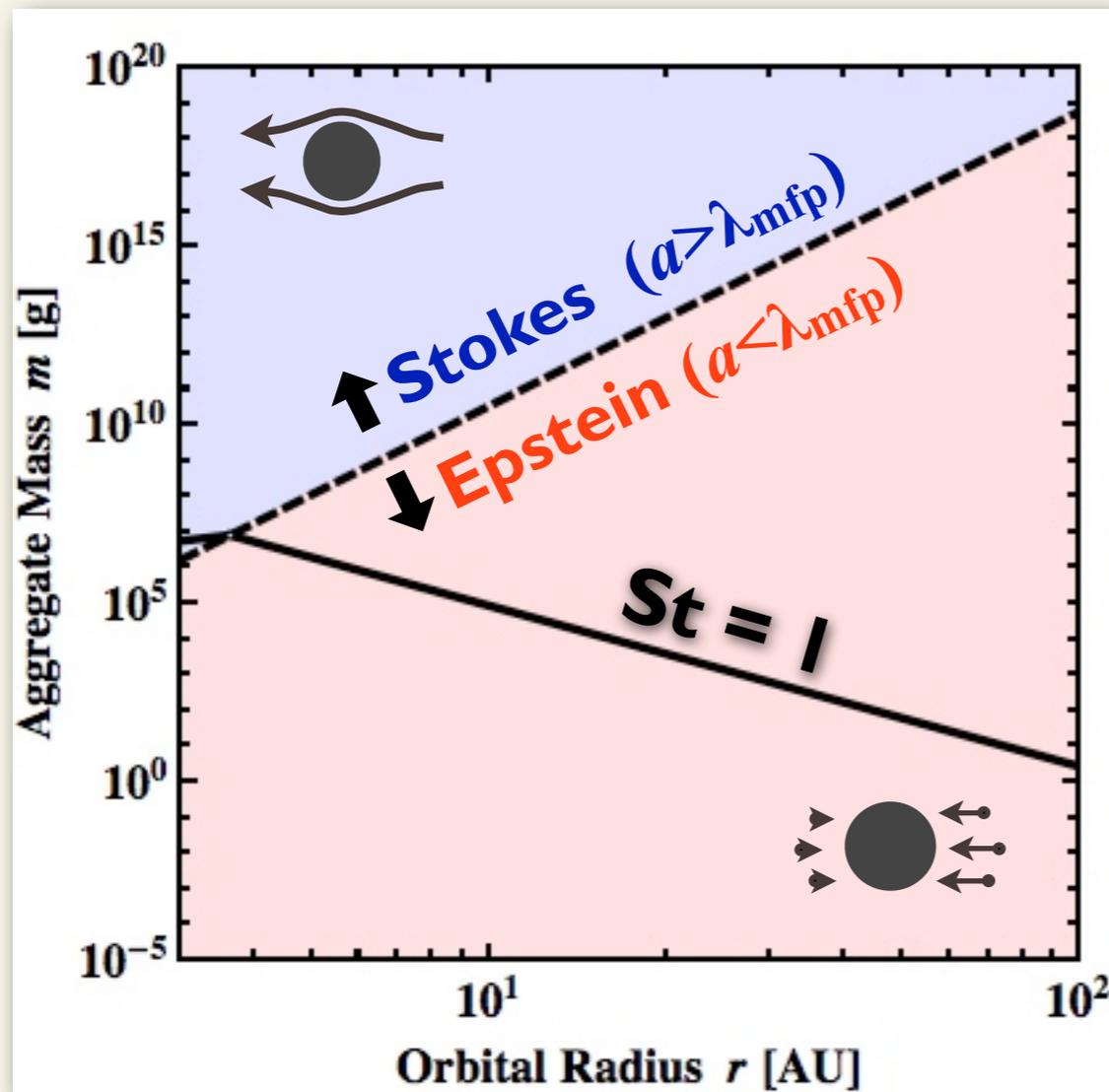
(Okuzumi et al. 2012)

Evolution of the Aggregate Internal Density

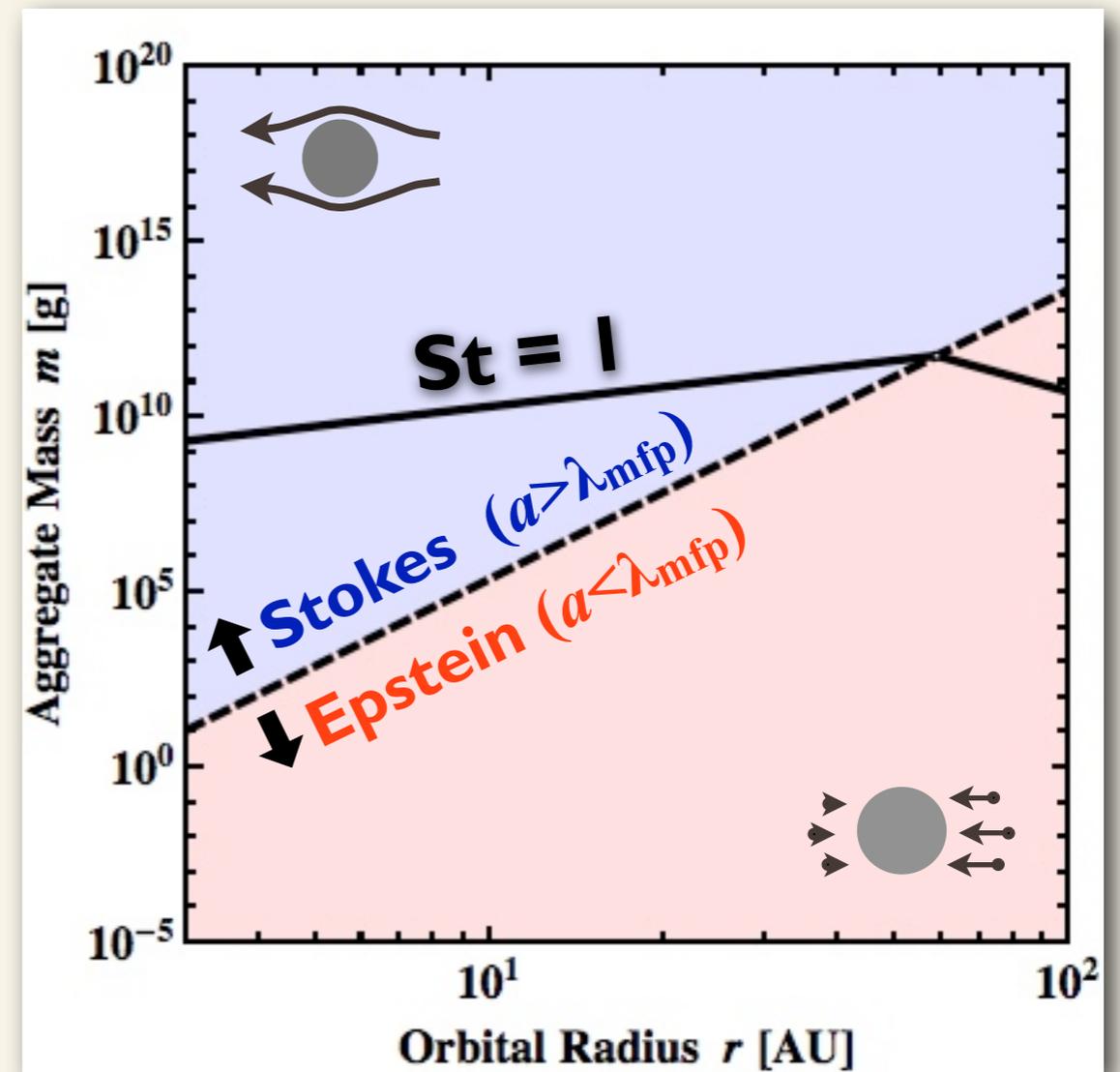


Higher Porosity \Rightarrow Wider Stokes Regime

Compact ($\phi=100\%$)



Fluffy ($\phi=10^{-5}$)



$St = 1$ in the **Epstein** regime

$St = 1$ in the **Stokes** regime

“Stokes Acceleration”

$$\text{Growth timescale } t_{\text{grow}} \equiv \left(\frac{d \ln m}{dt} \right)^{-1} = \frac{m}{\rho_d \sigma \Delta v}$$

● In the Epstein Regime ($a \ll \lambda_{\text{mfp}}$):

$$t_{\text{grow}} \sim 100 \left(\frac{\Sigma_g / \Sigma_d}{100} \right) \Omega^{-1}$$

(Takeuchi & Lin 05; Brauer et al. 08; Birnstiel et al. 12)

This is *not* short enough to overcome the drift barrier! (see Brauer et al. 08)

● However, in the Stokes Regime ($a \gg \lambda_{\text{mfp}}$):

$$t_{\text{grow}} \sim 100 \left(\frac{\lambda_{\text{mfp}}}{a} \right) \left(\frac{\Sigma_g / \Sigma_d}{100} \right) \Omega^{-1}$$

(Okuzumi et al. 12)

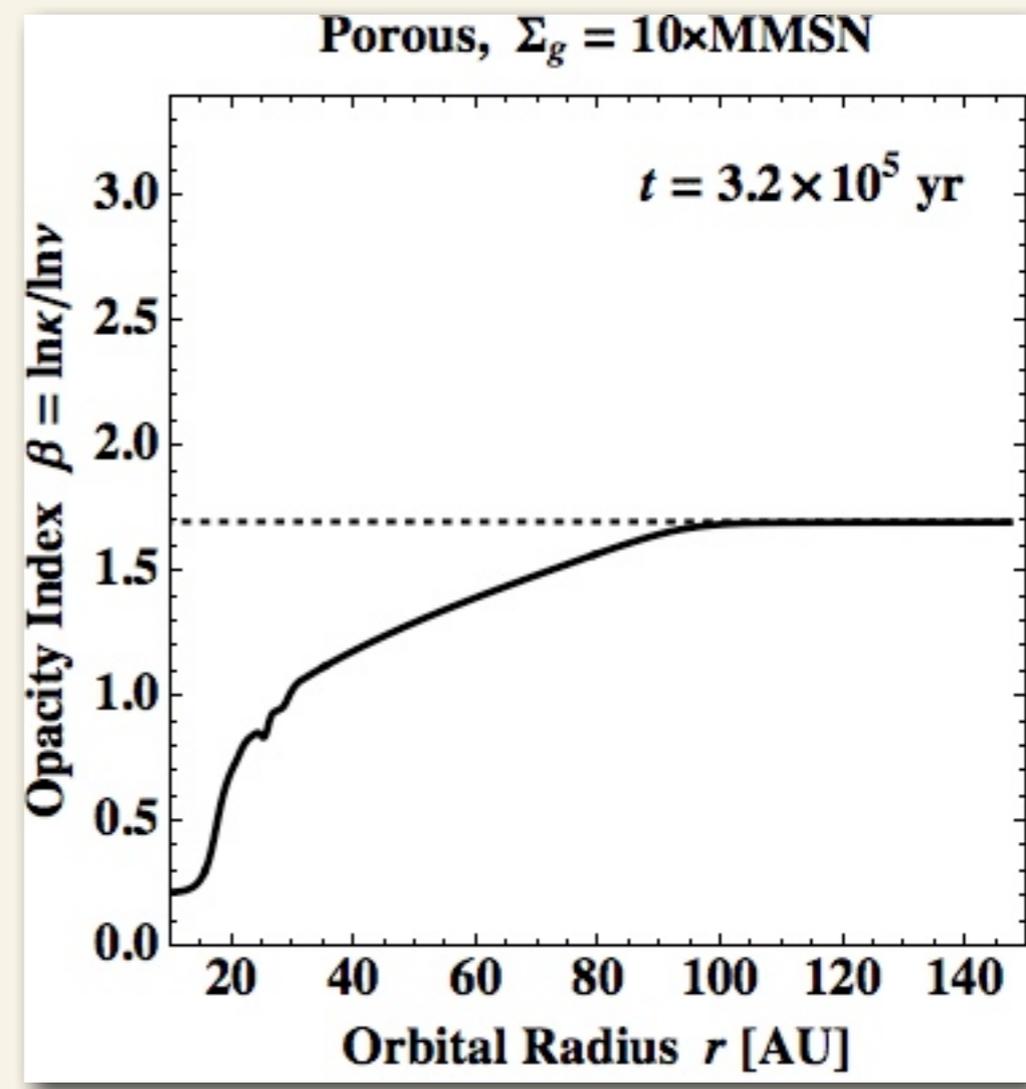
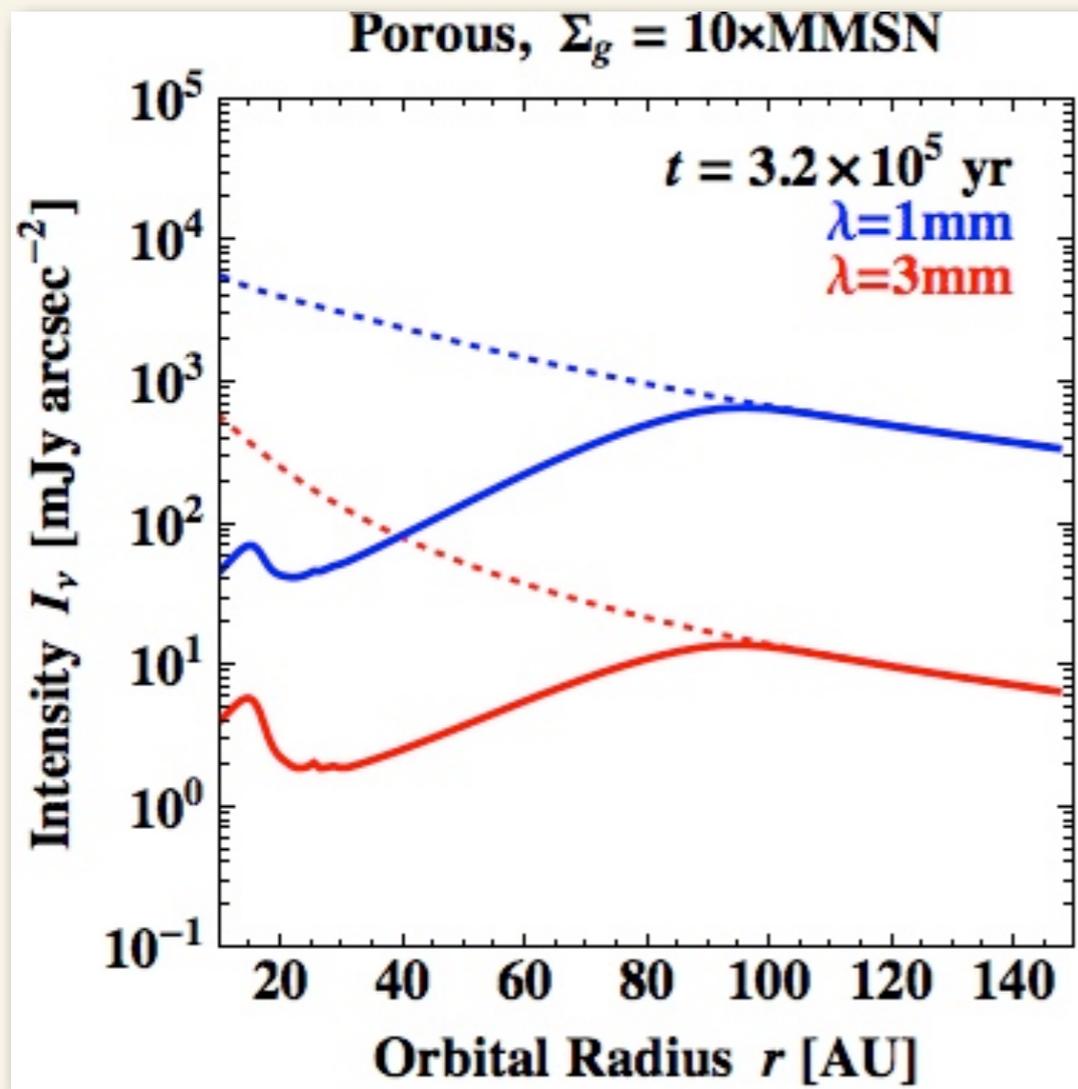
➔ **Rapid Growth in the “Deep” Stokes Regime ($a \gg \lambda_{\text{mfp}}$)**

(see also Fig. 11 of Birnstiel et al. '10 and Fig. 3 of Zsom et al. '11)

Millimeter Appearance of the Stokes Acceleration

(Preliminary!)

- $\Sigma_g = 10 \times \text{MMSN}$, $(\Sigma_d/\Sigma_g)_{\text{init}} = 0.01$
- κ_ν calculated with the Maxwell-Garnett effective medium theory

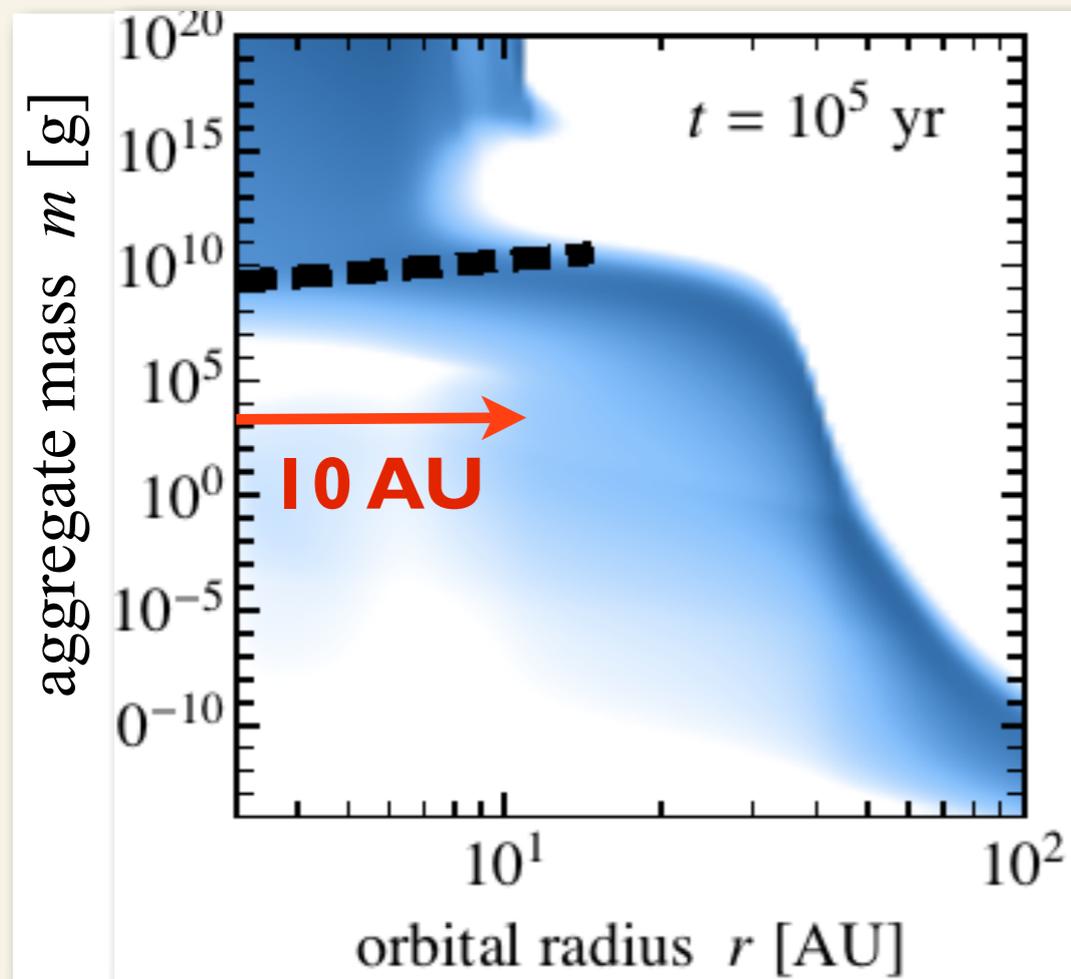


Summary

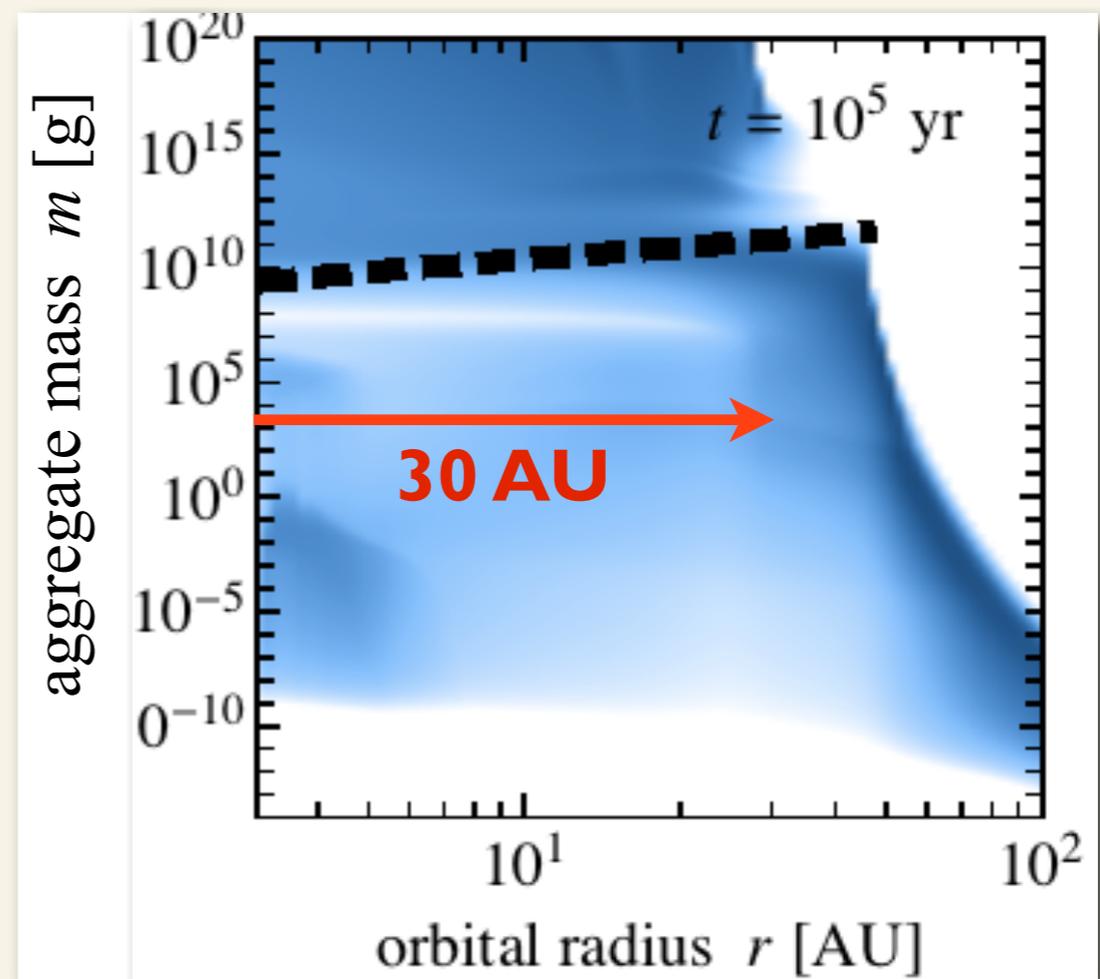
- Porous aggregates grow rapidly thanks to the fast decoupling in the Stokes drag regime (“**Stokes Acceleration**”).
- **Icy planetesimals can form through direct collisional growth at $\lesssim 10$ AU.**
- Caveats:
 - Rocky planetesimal formation is still a mystery! (Fragmentation barrier needs to be overcome.)
 - Non-collisional compression (by headwind / self gravity) needs to be taken into account.

Disk Mass Dependence

MMSN



10×MMSN



Higher $\Sigma g \rightarrow$ Shorter $\lambda_{\text{mfp}} \rightarrow$ Even Wider Stokes Region

